

# COMSOL used for simulating biological remodelling

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# Joint work with

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- Markus Michael Knodel

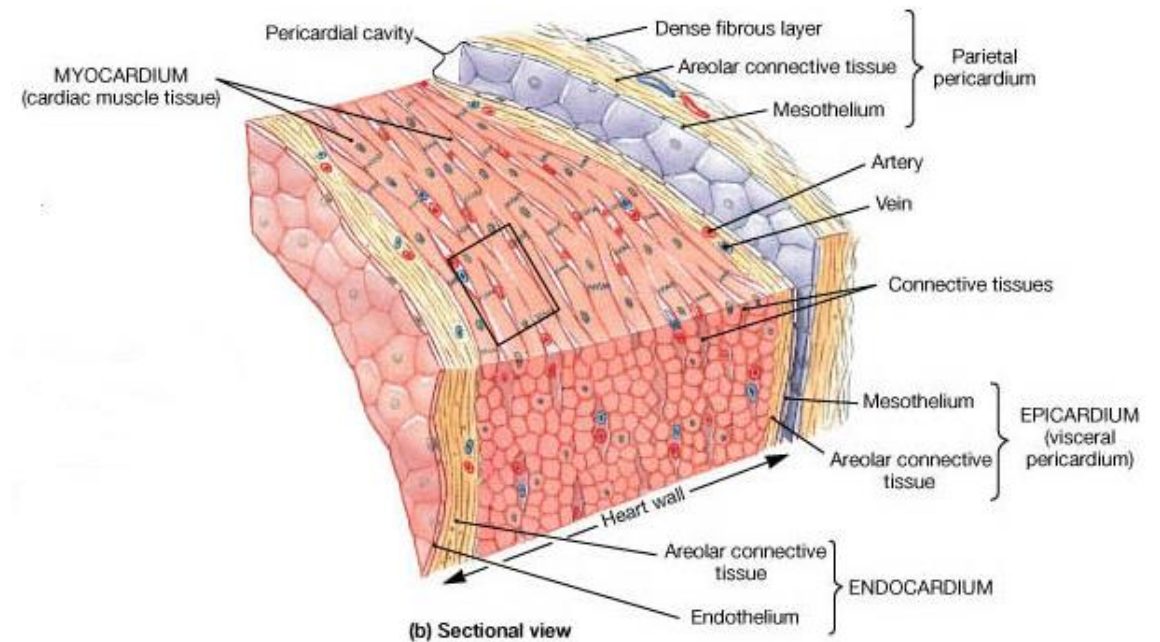
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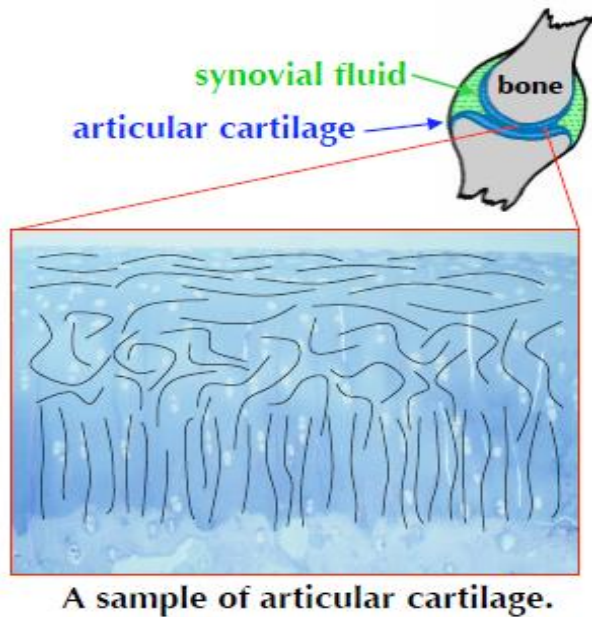
# Biological tissues: an overview

- Highly complex physical systems
- Hydrated, fibre-reinforced, heterogeneous and anisotropic porous media.
- Description of mechanical interactions
- Anelastic distortions



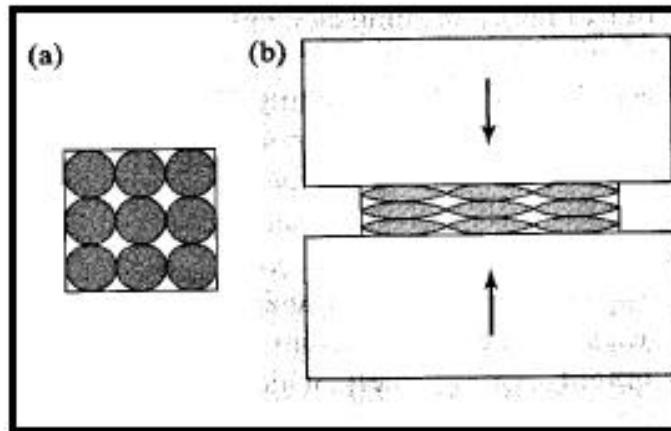
# Articular cartilage

In articular cartilage, the fibres are distributed in a non-uniform way and this influences, for example, the stiffness of the tissue and the flow of the interstitial fluid.

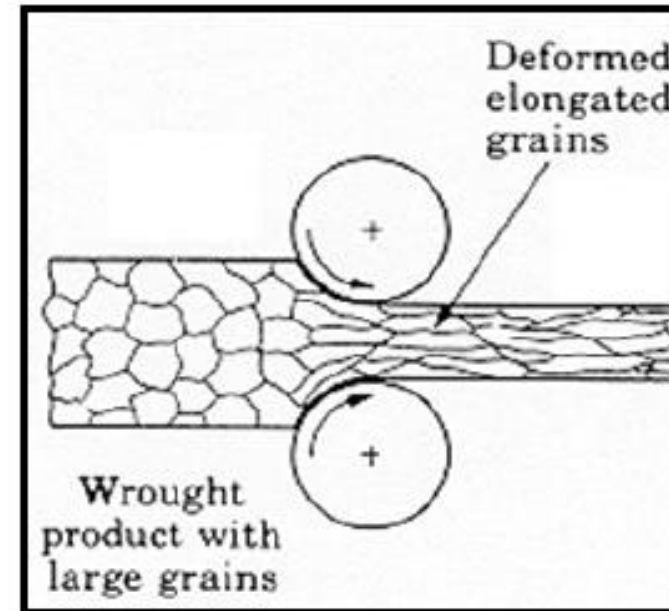


- Deep zone (oriented perpendicularly to the interface)
- Middle zone (random distributed)
- Upper zone (the fibres are parallel to the interface)
- *Transversely isotropic* behaviour

# Remodelling: change of internal structure



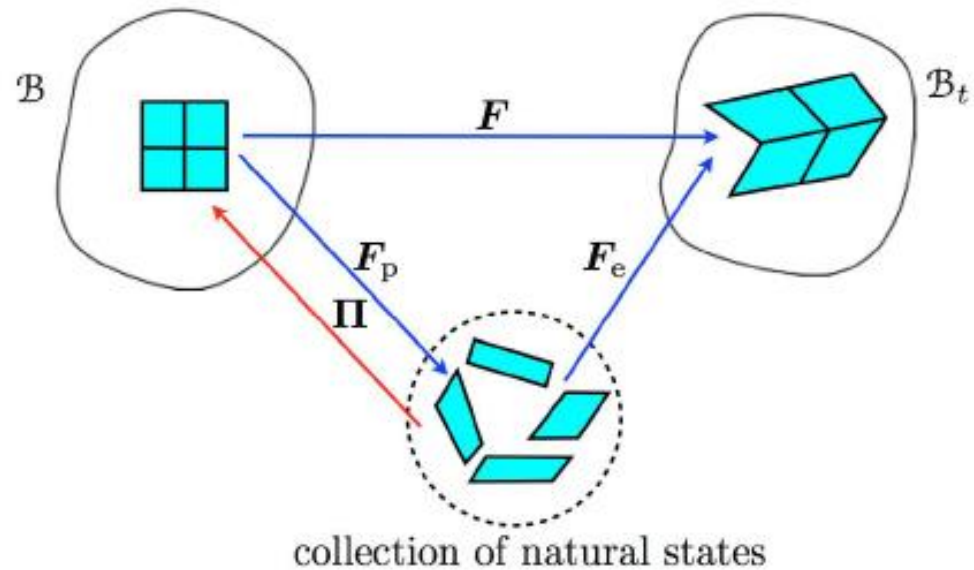
Unconfined compression



Cold rolling

The change of the body's shape (**visible phenomenon**) is accompanied by a reorganization of its internal structure (**hidden phenomenon**), which causes macroscopic variations of the mechanical properties of the material.

# The implant tensor



- Reference configuration  $\mathcal{B}$
- Actual configuration  $\mathcal{B}_t$
- Deformation gradient  $F$
- BKL decomposition:  $F = F_e F_p$
- Epstein-Maugin decomposition:  $F = F_e \Pi^{-1}$
- $\Pi$  is said to be the implant tensor

# Fibre pattern

- **Unit sphere:** set of all unit vectors emanating from  $X \in \mathcal{B}$

$$\mathbb{S}_X^2 \mathcal{B} = \{\mathbf{m}_X \in T_X \mathcal{B} : \|\mathbf{m}_X\| = 1\}$$

- **Probability density** that a fibre is aligned along  $\mathbf{m}_X$

$$\wp_X : \mathbb{S}_X^2 \mathcal{B} \rightarrow \mathbb{R}_0^+, \mathbf{m}_X \rightarrow \wp_X(\mathbf{m}_X)$$

- **Directional average** of a physical quantity associated with the fibres  $\mathfrak{F}_X : \mathbb{S}_X^2 \mathcal{B} \rightarrow \mathbb{R}$

$$\langle\langle \mathfrak{F}_X \rangle\rangle = \int_{\mathbb{S}_X^2 \mathcal{B}} \mathfrak{F}_X(\mathbf{m}_X) \wp_X(\mathbf{m}_X) = \int_0^{2\pi} \int_0^\pi \mathfrak{F}_X(\hat{\mathbf{m}}_X(\Theta, \Phi)) \wp_X(\hat{\mathbf{m}}_X(\Theta, \Phi)) \sin(\Theta) d\Theta d\Phi$$

where

$$\mathbf{m}_X = \hat{\mathbf{m}}_X(\Theta, \Phi) = \sin \Theta \cos \Phi \mathbf{e}_1 + \sin \Theta \sin \Phi \mathbf{e}_2 + \cos \Theta \mathbf{e}_3$$

with  $(\Theta, \Phi) \in [0, \pi[ \times [0, 2\pi]$ , and  $\{\mathbf{e}_\alpha\}_{\alpha=1}^3$  orthonormal basis in  $T_X \mathcal{B}$ .

- **Transverse isotropy**

$$\exists \mathbf{m}_0 \text{ such that, for } \mathbf{H}_0 : \mathbf{H}_0 \mathbf{m}_0 = \pm \mathbf{m}_0, \wp_X(\mathbf{H}_0 \mathbf{m}_X) = \wp_X(\pm \mathbf{m}_X) \forall X \in \mathcal{B}$$

- **Parity Symmetry**

$$\wp_X(\mathbf{m}_X) = \wp_X(-\mathbf{m}_X), \text{ for all } X \in \mathcal{B}$$

## Constitutive framework: strain energy function

- Hyperelastic behaviour from the collection of natural states
- Strain energy function

$$\begin{aligned}W_{\text{R}}(\mathbf{C}, X, t) &= \frac{1}{J_{\Pi}(X, t)} \hat{W}_{\kappa}(\mathbf{C}_{\text{e}}(X, t)) \\ &= \frac{1}{J_{\Pi}(X, t)} (\Phi_{\text{sv}} \hat{U}(J_{\text{e}}) + \Phi_{\text{0sv}} \hat{W}_0(\mathbf{C}_{\text{e}}) + \Phi_{\text{1sv}} \hat{W}_{\text{e}}(\mathbf{C}_{\text{e}}))\end{aligned}$$

with

- $J_{\Pi}(X, t) = \det(\Pi(X, t))$ ,  $J_{\text{e}}(X, t) = \det(\mathbf{F}_{\text{e}}(X, t))$ ;
- $\Phi_{\text{sv}}(X, t) = J_{\text{e}}(X, t)\phi_{\text{s}}(X)$ ,  $\Phi_{\text{0sv}}(X, t) = J_{\text{e}}(X, t)\phi_{\text{0sv}}(X)$ ,  $\Phi_{\text{1sv}}(X, t) = J_{\text{e}}(X, t)\phi_{\text{1sv}}(X)$  are, respectively, the volumetric fractions of the solid phase, the matrix and the fibres, such that  $\phi_{\text{0s}} + \phi_{\text{1s}} = \phi_{\text{s}}$ ;
- $\mathbf{C}_{\text{e}} = \mathbf{F}_{\text{e}}^T \cdot \mathbf{F}_{\text{e}} = \Pi^T \mathbf{C} \Pi$  is the elastic part of the right Cauchy-Green deformation tensor  $\mathbf{C}$ .



### Strain energy of the matrix

$$\hat{W}_0(\mathbf{C}_e) = \alpha_0 \frac{\exp(\alpha_1[I_1 - 3] + \alpha_2[I_2 - 3])}{[I_3]^{\alpha_3}}$$

$$I_1 = \text{tr}(\mathbf{C}_e)$$

$$I_2 = \frac{1}{2}[(\text{tr}(\mathbf{C}_e))^2 - \text{tr}(\mathbf{C}_e^2)]$$

$$I_3 = \det(\mathbf{C}_e)$$

$$\alpha_1 + 2\alpha_2 = \alpha_3 = 1$$

### Ensemble potential

$$\hat{W}_e(\mathbf{C}_e) = \hat{W}_{li}(\mathbf{C}_e) + \langle\langle \hat{W}_{la}(\mathbf{C}_e, \mathbf{m}) \rangle\rangle$$

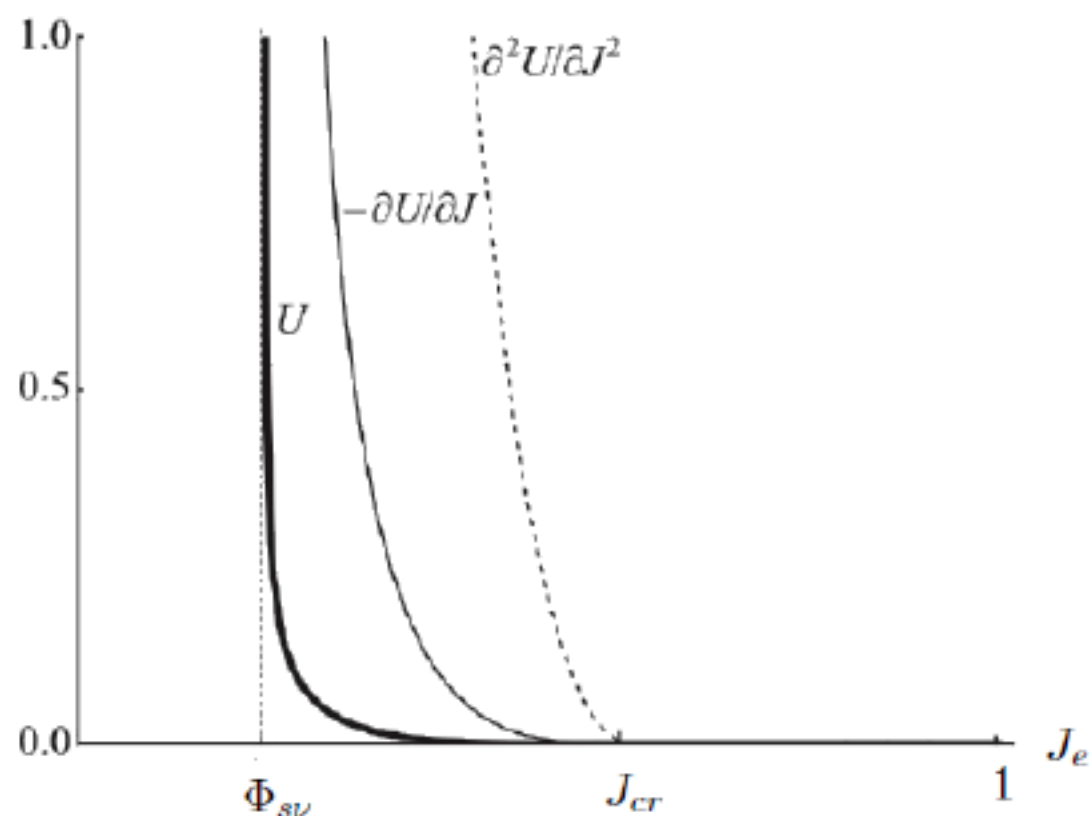
$$\hat{W}_{li} = \hat{W}_0(\mathbf{C}_e)$$

$$\hat{W}_{la}(\mathbf{C}_e, \mathbf{m}) = \mathcal{H}(I_{4e} - 1) \frac{1}{2} c [I_{4e} - 1]^2$$

$$I_{4e} = \mathbf{C}_e : \mathbf{a}, \quad \mathbf{a} = \mathbf{m} \otimes \mathbf{m}$$

### Penalty term

$$\hat{U}(J_e) = \mathcal{H}(J_{cr} - J_e)(J_e - J_{cr})^{2q}(J - \Phi_{sv})^{-r}$$



$$J_{crit} \in (\phi_{sv}, 1], q \geq 2 \in \mathbb{N}, r \in (0, 1]$$

# Constitutive framework: Darcy's law

Linear relation between the filtration velocity and the gradient of pressure

- spatial formulation:  $\mathbf{q} = \phi_f(\mathbf{v}_f - \mathbf{v}_s) = \mathbf{k} \text{grad } p$
- material formulation:  $\mathbf{Q} = J\mathbf{F}^{-1}\mathbf{q} = \mathbf{K} \text{Grad } p$

where

- $\mathbf{v}_f$  is the velocity of the fluid phase and  $\mathbf{v}_s$  is the velocity of the solid phase;
- $\phi_f = 1 - \phi_s$  is the volumetric fraction of the fluid phase;
- $\mathbf{k}$  is the spatial permeability tensor and  $\mathbf{K}$  is the material permeability tensor.

In particular

$$\mathbf{K} = \hat{\mathbf{K}}(\mathbf{F}, \mathbf{\Pi}, \zeta) = J \hat{k}_0(J, J_{\mathbf{\Pi}}, \zeta) \mathbf{C}^{-1} + J^{-1} \hat{k}_0(J, J_{\mathbf{\Pi}}, \zeta) \mathbf{\Pi} \left\langle \left\langle \frac{\mathbf{a}}{I_{4e}} \right\rangle \right\rangle \mathbf{\Pi}^T$$

where  $\zeta$  is an axial coordinate and  $k_0 = \hat{k}_0(J, J_{\mathbf{\Pi}}, \zeta)$  is a Holmes-Mow type scalar permeability.

# Problem setting (1/2)

The unknowns of the problem are the motion  $\chi$ , the pore pressure  $p$  and the implant tensor  $\mathbf{\Pi}$ .  
The system of equations to solve is

$$\begin{aligned} \dot{J} - \text{Div}[\mathbf{K} \text{Grad } p] &= 0, \\ \text{Div}(-J p \mathbf{g}^{-1} \mathbf{F}^{-T} + \mathbf{P}_{sc}) &= \mathbf{0}, \\ \text{sym}(\mathbf{\Lambda} \mathbf{C}^{-1}) &= \xi_p \left[ \mathbf{S} - \frac{1}{3} \text{tr}(\mathbf{C} \mathbf{S}) \mathbf{C}^{-1} \right]. \end{aligned}$$

In addition, we suppose  $J_{\mathbf{\Pi}} = 1$  and  $\xi_p = \lambda_0 \phi_s^2 \frac{[\|\text{dev}(\boldsymbol{\sigma})\| - \sqrt{(2/3)} \sigma_y]_+}{\|\text{dev}(\boldsymbol{\sigma})\|}$ , with  $\boldsymbol{\sigma}$  being the Cauchy stress tensor,  $\sigma_y$  is a yield stress and  $[\cdot]_+$  extracts the positive part of the function to which it is applied

## Problem setting (2/2)

- The first equation is the balance of mass.
- The second equation is the moment balance equation, where  $\mathbf{P}_{sc} = \hat{\mathbf{P}}_{sc}(\mathbf{F}, \mathbf{\Pi})$  is the constitutive part of the first Piola-Kirchoff stress tensor.
- The third equation is the law of evolution for the implant tensor  $\mathbf{\Pi}$ , where  $\mathbf{\Lambda} = \dot{\mathbf{\Pi}}\mathbf{\Pi}^{-1}$  is the tensor of *inomegeneities velocity*.
- A pseudo-Gaussian probability distribution has been chosen

$$\wp(\Theta) = \hat{\wp}(\Theta, X, t) = \frac{\hat{\gamma}(\Theta, X, t)}{\int_0^{2\pi} \hat{\gamma}(\Theta', X, t) \sin(\Theta') d\Theta'}$$
$$\gamma(\Theta) = \hat{\gamma}(\Theta, X, t) = \exp\left(-\frac{[\Theta - Q(X)]^2}{2\omega(X)}\right),$$

where  $Q$  is the mean angle and  $\omega$  is the variance.

# Closing the system

Let us introduce the polar decomposition of the implant tensor

$$\mathbf{\Pi} = \mathbf{H}.\mathbf{R} = \mathbf{H}\mathbf{G}\mathbf{R}$$

where

$\mathbf{H}$  is a symmetric positive definite tensor;

$\mathbf{R}$  is a rotation tensor;

By imposing  $R^\beta_\alpha = \delta^\beta_\alpha$ , we obtain

$$\mathbf{\Pi} = \mathbf{H}\mathbf{G}$$

and

$$\mathbf{\Lambda} = \dot{\mathbf{\Pi}}\mathbf{\Pi}^{-1} = \dot{\mathbf{H}}\mathbf{H}^{-1}.$$

We can solve the flow rule for a symmetric tensor.

# Simulations (1/2)

We simulated an unconfined compression test for a cylindrical specimen. In this case we imposed the following boundary conditions

$$\begin{aligned}\chi(X, t) &= \chi(X, 0) = X, \\ (\mathbf{K} \text{ Grad } p) \cdot \mathbf{N} &= 0\end{aligned}$$

on the lower boundary  $\Gamma_l$ ,

$$\begin{aligned}(-J p \mathbf{g}^{-1} \mathbf{F}^{-T} + \mathbf{P}_{sc}) \cdot \mathbf{N} &= \mathbf{0}, \\ p &= 0\end{aligned}$$

on the lateral boundary  $\Gamma_L$  and

$$\begin{aligned}\chi^z(X, t) &= g(t), \\ \mathbf{K} (\text{Grad } p) \cdot \mathbf{N} &= 0\end{aligned}$$

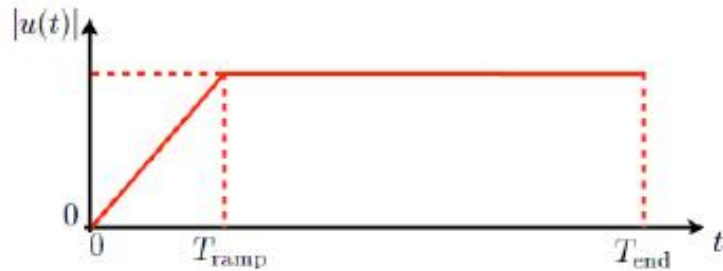
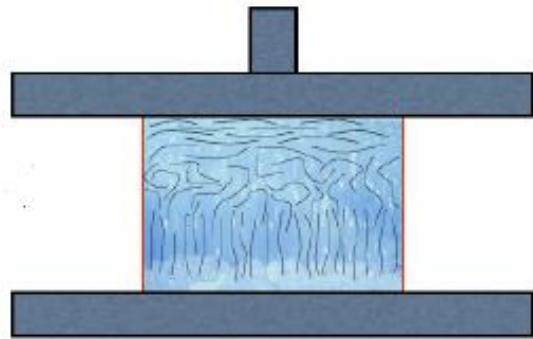
on the upper boundary  $\Gamma_u$ .

# Simulations (2/2)

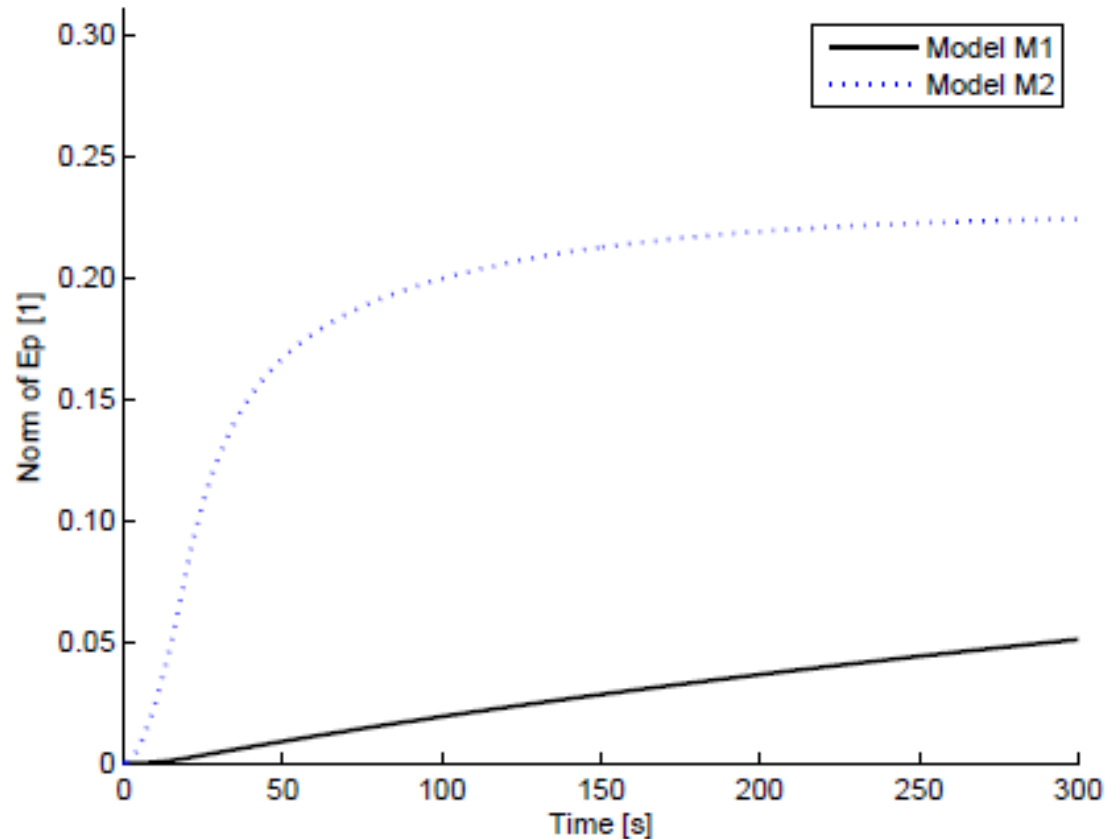
In the previous equations,  $g$  is an imposed displacement given by

$$g = \begin{cases} L - \frac{t}{T_{\text{ramp}}} u_T, & t \in [0, T_{\text{ramp}}] \\ L - u_T, & t \in [T_{\text{ramp}}, T_{\text{end}}] \end{cases}$$

where  $T_{\text{ramp}}$  is the final instant of time of the loading ramp and  $u_T$  is a reference displacement.



# Results: plastic strain behaviour



$E_p = \frac{1}{2} [\mathbf{\Pi}^{-1} \cdot \mathbf{\Pi} - \mathbf{G}]$  is the Almansi-Euler like strain tensor associated to the anelastic distortions



## Conclusions and future works

- Study of the mechanical properties of biological tissues
- Anelastic distortion
- Possibility to implement other flow rules: take into account correlation of the fibres
- Coupling anelastic distortions and evolution of the fibre pattern
- Introduction of the Forchheimer correction for the fluid flow

Thank you for your attention!